

(from: An Introduction to Optoelectronics
Alan Rogers)

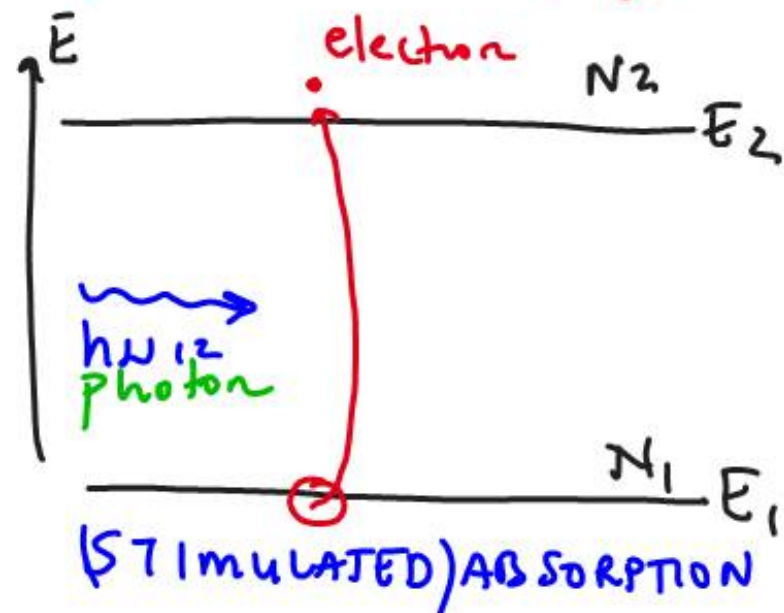
23.10.2012

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QUANTITATIVE EVALUATION

1. 2-level Systems

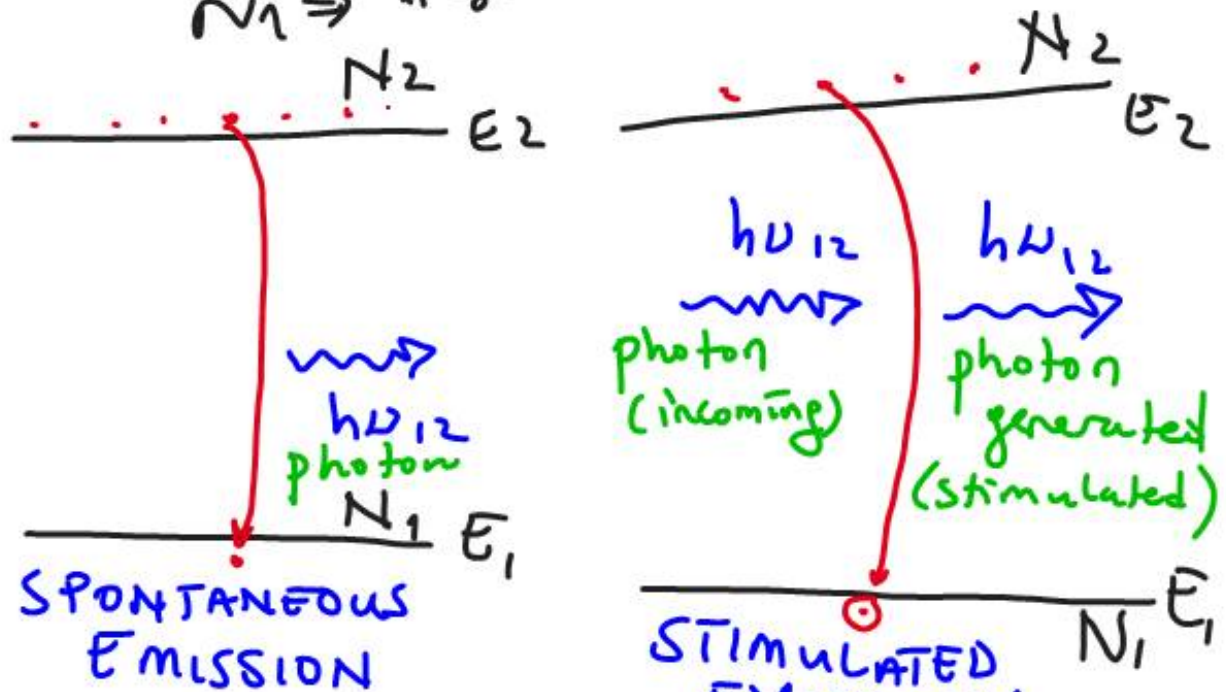
$N_2 \Rightarrow$ # of electrons at E_2
 $N_1 \Rightarrow$ # of " at E_1



The rate of this process

$$R_{12} = B_{12} \cdot \rho_{\nu} \cdot N_1$$

ρ_{ν} : density of photons (per volume, per frequency) at energy $h\nu_{12}$



$$S_{21} = A_{21} \cdot N_2$$

$$R_{21} = B_{21} \cdot \rho_{\nu} \cdot N_2$$

B_{12}, A_{21}, B_{21} : Einstein's coefficients

In equilibrium upward transitions = downward transitions

$$R_{12} = S_{21} + R_{21}$$

$$B_{12} \rho_{\nu} N_1 = A_{21} N_2 + B_{21} \rho_{\nu} N_2$$

$$B_{12} \rho_{\nu} = A_{21} \frac{N_2}{N_1} + B_{21} \rho_{\nu} \frac{N_2}{N_1}$$

$$\rho_{\nu} \left(B_{12} - B_{21} \frac{N_2}{N_1} \right) = A_{21} \frac{N_2}{N_1}$$

$$\rho_{\nu} = \frac{A_{21} N_2 / N_1}{B_{12} - B_{21} \frac{N_2}{N_1}} = \frac{A_{21} \frac{N_2}{N_1}}{\frac{B_{12} N_1 - B_{21} N_2}{N_1}} = \frac{A_{21}}{B_{12} \frac{N_1}{N_2} - B_{21}} = \frac{A_{21} / B_{21}}{\frac{B_{12} N_1}{B_{21} N_2} - 1}$$

$$\rho_{\nu} = \frac{A_{21} / B_{21}}{\left(\frac{B_{12} N_1}{B_{21} N_2} \right) - 1}$$

$$\frac{N_1}{N_2} = \exp\left(-\frac{E_1 - E_2}{kT}\right)$$

$$E_2 - E_1 = h\nu_{12}$$

$$\rho_\nu = \frac{A_{21}/B_{21}}{\left(\frac{B_{12}}{B_{21}}\right)\left(\exp\frac{h\nu_{12}}{kT} - 1\right)}$$

For Blackbody radiation at equilibrium

$$\rho_\nu = \frac{8\pi h\nu^3}{c^3} \frac{1}{\left(\exp\frac{h\nu}{kT} - 1\right)}$$

$$B_{12} = B_{21}$$

$$A_{21} = B_{21} \cdot \frac{8\pi h\nu^3}{c^3}$$

**EINSTEIN'S
RELATIONS**

Results:

1) The ratio of stimulated to spontaneous emission

$$S = \frac{R_{21}}{S_{21}} = \frac{\rho_\nu \cdot B_{21} \cdot \cancel{N_2}}{A_{21} \cdot \cancel{N_2}}$$

$$S = \rho_\nu \cdot \frac{B_{21}}{A_{21}}$$

$$S = \rho_\nu \cdot \frac{c^3}{8\pi h \nu^3}$$

$$S = \frac{8\pi h \nu^3}{c^3} \cdot \frac{c^3}{8\pi h \nu^3} \cdot \frac{1}{\left(\exp \frac{h\nu}{kT} - 1\right)}$$

$$S = \frac{1}{\left(\exp \frac{h\nu}{kT} - 1\right)}$$

$$S = \frac{1}{\exp \frac{h\nu}{kT} - 1}$$

for He-Ne discharge at $T = 370 \text{ K}$

$$\lambda = 632.8 \text{ nm}$$

$$\nu = 4.74 \times 10^{14} \text{ Hz}$$

$$S = 2 \times 10^{-27}$$

Stimulated Emission is very unlikely to occur in a 2-level system in equilibrium.

Mostly spontaneous emission occurs.

$$2) \quad B_{21} = A_{21} \cdot \frac{c^3}{8\pi h \nu^3}$$

$$\frac{B_{21}}{A_{21}} = \frac{c^3}{8\pi h \nu^3}$$

Stimulated ← B_{21}
Spontaneous ← A_{21}

$$\frac{B_{21}}{A_{21}} \propto \nu^{-3}$$

This means it is harder to obtain stimulated emission at higher frequencies.

IR Laser → Easier to make
Blue Laser → Difficult to make

Light Amplification by Stimulated Emission of Radiation

LASER

Microwave Amplification " " " " " "

MASER

First MASERS were invented. } Since it is easier to make
Then LASERS " " } MASERS than LASERS

$$(\nu_{\text{maser}} < \nu_{\text{laser}})$$

Radio frequency
 $\sim 10^9 \text{ Hz}$

Light frequency
 $\sim 10^{14} \text{ Hz}$

In order to have lasing action we have to disturb the equilibrium and make $\frac{N_2}{N_1} > 1$, and to increase ρ_ω as well:

$$N_2 \rho_\omega B_{21} \gg N_1 A_{21}$$

$$\rho_\omega \gg \frac{A_{21}}{B_{21}}$$

$$\rho_\omega \gg \frac{8\pi h \omega^3}{c^3}$$

However, increasing ρ_ω also increases the absorption when ρ_ω gets large

$$N_1 \cdot \rho_\omega B_{12} \cong N_2 \rho_\omega B_{21}$$

ABSORPTION
STIMULATED EMISSION

for $B_{12} \approx B_{21}$

$$N_1 = N_2$$

In 2-Level system:

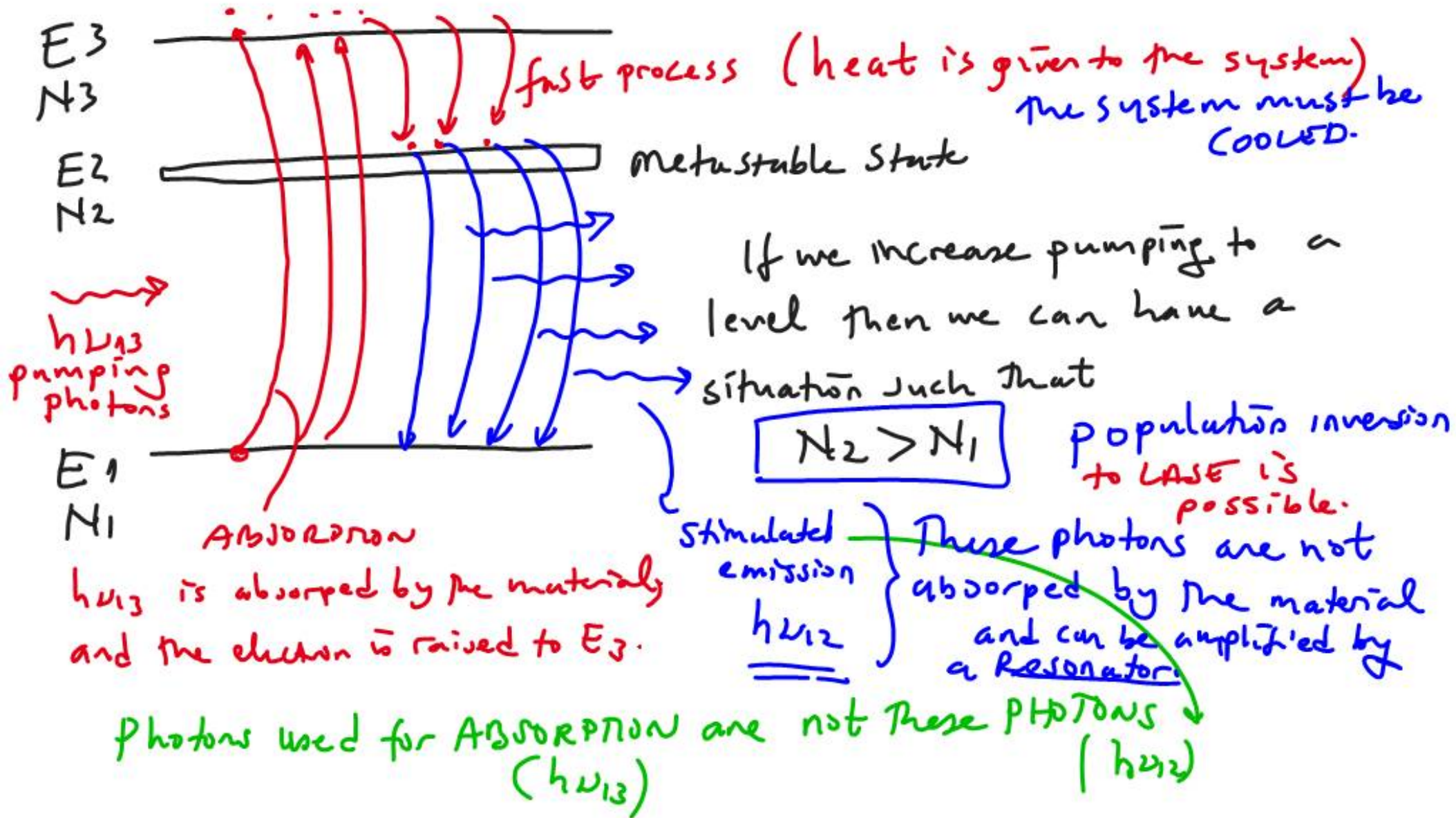
At most N_2 can be equal to N_1 ,
but cannot be bigger than N_1 .

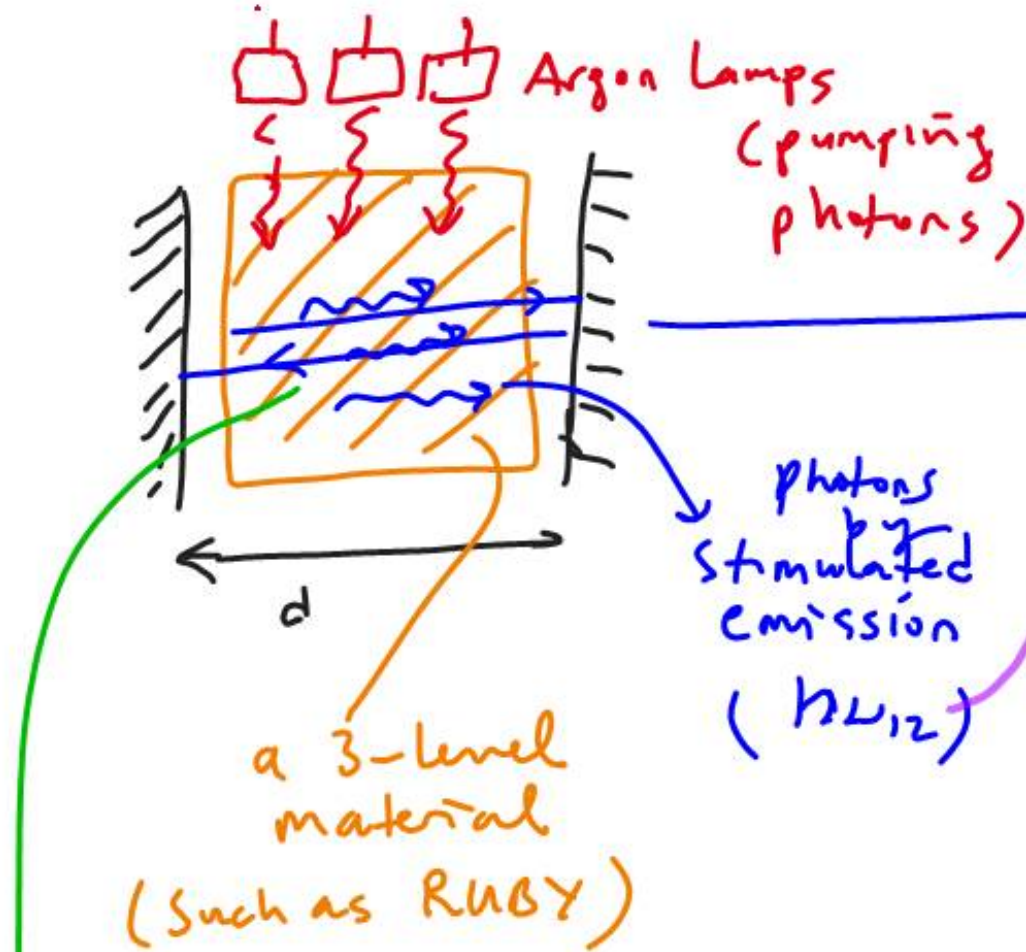
Hence, we can not increase N_2 higher than N_1
by increasing the value of the photon field (ρ_{ν})
by pumping the light (photons). Because, these photons
are used for absorption as well

POPULATION INVERSION can not be obtained.

Therefore it cannot LASE.

2) A 3-LEVEL SYSTEM





Argon Lamps
(pumping photons) $h\nu_{13}$

These photons are not absorbed again ($h\nu_{12} < h\nu_{13}$)

Laser beam

Only pulse operation is possible in 3-level LASER:

When all electrons are used in E_2 to emit stimulated photons
Then population inversion is no longer sustained.

Therefore a pulse of laser is obtained by the stimulated emission photons. The other pulse can only be generated when pumping and absorption populates E_2 level of energy with enough electrons such that N_2 becomes greater than N_1

a 3-level material
(Such as RUBY)

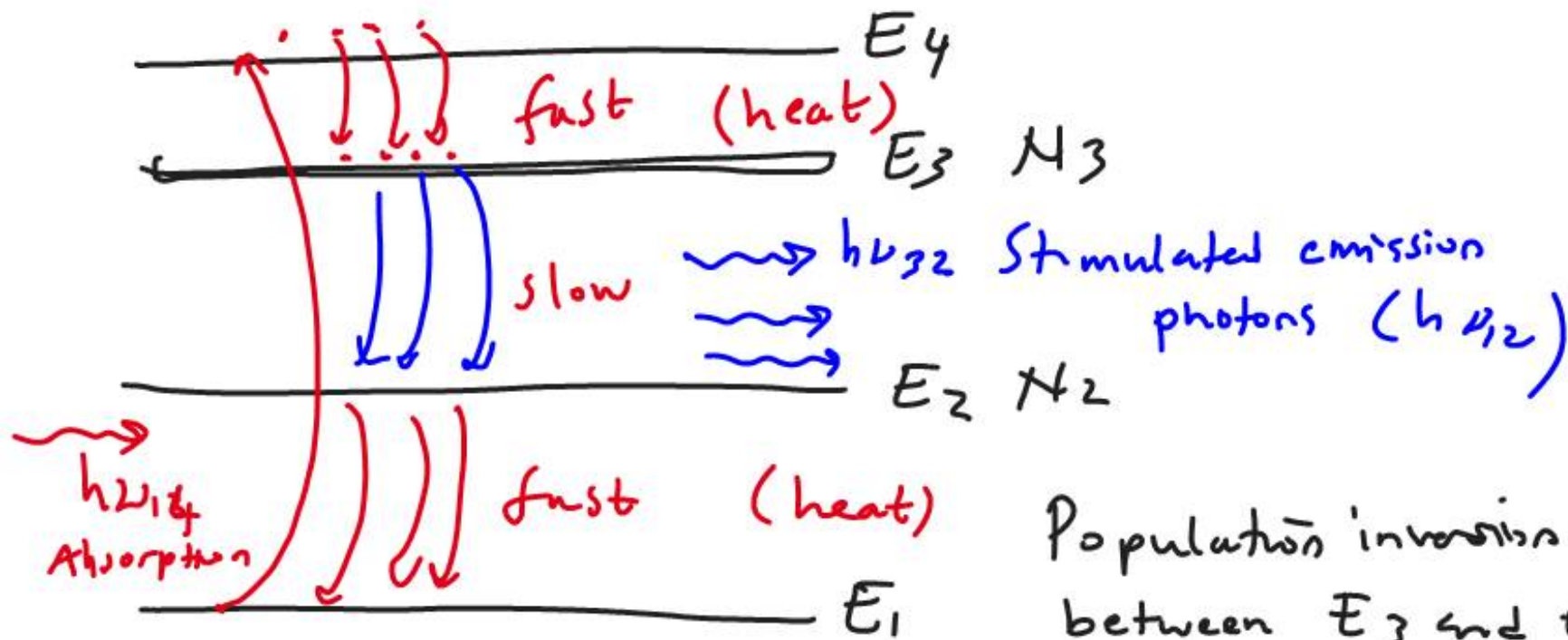
photons stimulated emission
($h\nu_{12}$)

stimulated photons are amplified by the resonator

i.e., population inversion is obtained again

$N_2 > N_1$

3- A 4-level Laser



Nd: Glass Laser

He-Ne Laser

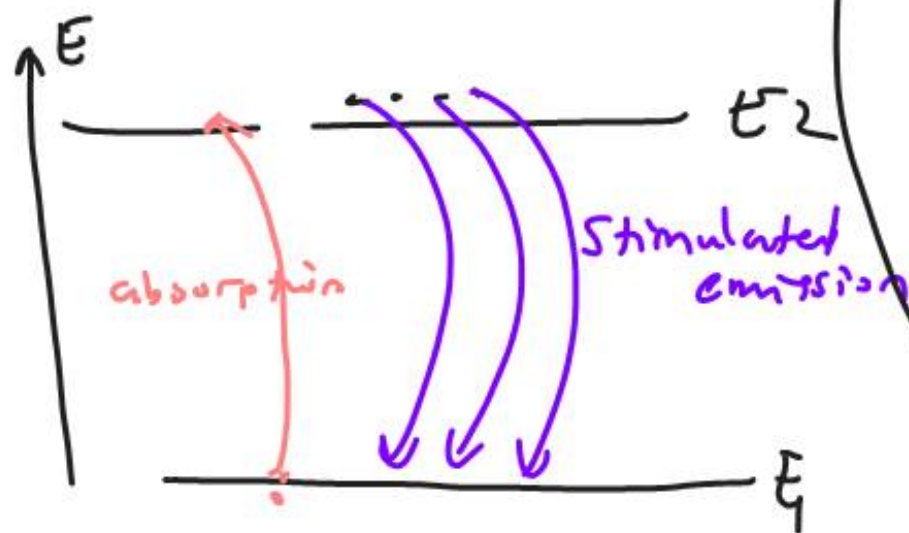
Population inversion takes place between E_3 and E_2 , ($N_3 > N_2$) and can be continuously, and easily Sustained.

Therefore, a continuous wave (CW) operation is possible in these Lasers.

if we have a Lasing system between energy levels of E_2 and E_1 :

The # of photons per unit volume (stimulated ^{emission} and _{absorption}) that are entering the

system is ρ_{12} .



$$\frac{d\rho_{12}}{dt} = \underbrace{N_2 \rho_{12} B_{21}}_{\text{stimulated emission}} - \underbrace{N_1 \rho_{12} B_{12}}_{\text{Absorption}}$$

$$\rho_{12} = \rho_{12} (h\nu_{12})$$

each photon has this energy

$$\rho_{12} = \frac{\rho_{12}}{h\nu_{12}}$$

$$\frac{1}{h\nu_{12}} \frac{d\rho_{12}}{dt} = \rho_{12} (N_2 B_{21} - N_1 B_{12})$$

$$\text{If } \beta_{12} = \beta_{21} = \beta$$

$$\frac{1}{h\nu_{12}} \frac{d\rho_{12}}{dt} = \rho_{12} \beta (N_2 - N_1)$$

$$\frac{d\rho_{12}}{dt} = \rho_{12} \beta \cdot h\nu_{12} (N_2 - N_1)$$

$$\int \frac{\frac{d\rho_{12}}{dt}}{\rho_{12}} = \int \beta \cdot h\nu_{12} (N_2 - N_1) dt \qquad \int \frac{d\rho_{12}}{\rho_{12}} = \int \beta \cdot h\nu_{12} (N_2 - N_1) dt$$

$$\ln \rho_{12} - \ln \rho_{120} = \beta \cdot h\nu_{12} (N_2 - N_1) \cdot t$$

$$\ln\left(\frac{\rho_{12}}{\rho_{120}}\right) = \beta \cdot h\nu_{12} (N_2 - N_1) t$$

$$\rho_{12} = \rho_{120} \exp[\beta h\nu_{12} (N_2 - N_1) t]$$



$$\rho_{\nu_2} = \rho_{\nu_0} \exp [B \cdot h\nu_2 \cdot (N_2 - N_1) t]$$

as $N_2 > N_1$ ρ_{ν_2} increases.

if this photon field (wave) moves with the speed of light (c)

then it travels $s = ct$ in time t . $t = \frac{s}{c}$

$$\rho_{\nu} = \rho_{\nu_0} \exp \left[\frac{B h \nu}{c} \cdot (N_2 - N_1) s \right]$$

Since the intensity of the field is proportional to P_{av}

$$I \propto P_{av}$$

$$I = I_0 \cdot \exp(\beta \cdot x)$$

↓ ↓
amplification factor distance (s)

amplification factor \Rightarrow GAIN COEFFICIENT

$$\beta \propto \frac{h\nu}{c} \cdot B \cdot (N_2 - N_1)$$

$$\beta \propto (N_2 - N_1)$$

* Laser is actually
an AMPLIFIER. It amplifies
The photon field (LIGHT)
Light Amplification by Stimulated
Emission of Radiation (LASER)

If we find a way to feedback the signal properly we can obtain an oscillator from an amplifier

The feedback can be obtained by using the Resonator and the system continuously produces a beam of identical photons

- Photons of
- same energy
 - equal phase
 - same direction
 -

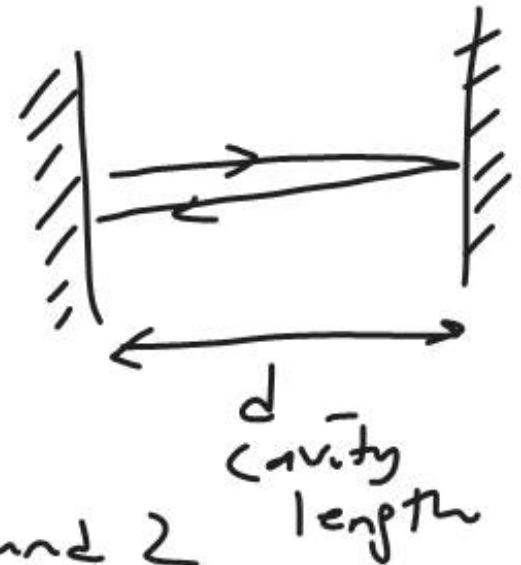
A LASER is actually the oscillator of a light of constant frequency, equal phase, and collimated beam of light (coherent)

$$\frac{I_f}{I_i} = R_1 R_2 \exp [(\beta - \alpha) 2d] > 1$$

$$\beta \propto \frac{h\nu}{c} \cdot B (N_2 - N_1)$$

I_f : Final intensity of the radiation

I_i : Initial intensity of the radiation



R_1, R_2 : reflectivities of the mirrors 1 and 2

$2d$: round trip distance in the resonator.
(d : cavity length)

α : loss per distance (unit length) in the medium

- Narrow band in wavelength spectrum (only 1 frequency is amplified)
- The output direction is normal to the mirrors and therefore highly collimated.
- All photons are locked in phase, and therefore, we have a coherent light within the limitations of the linewidth of the transitions ($\Delta\nu$)
- The light can be very intense

The LASER LIGHT:

- pure in wavelength and phase
- intense
- well collimated

and therefore it can be easily controlled and modulated.